

First-order nature of the ferromagnetic phase transition in (La-Ca)MnO₃ near optimal dopingC. P. Adams,^{1,2,*} J. W. Lynn,^{1,2} V. N. Smolyaninova,^{2,†} A. Biswas,² R. L. Greene,² W. Ratcliff II,³
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Neutron scattering has been used to study the nature of the ferromagnetic transition in single crystals of La_{0.7}Ca_{0.3}MnO₃ and La_{0.8}Ca_{0.2}MnO₃, and polycrystalline samples of La_{0.67}Ca_{0.33}MnO₃ and La_{5/8}Ca_{3/8}MnO₃ where the naturally occurring ¹⁶O can be replaced with the ¹⁸O isotope. Small angle neutron scattering on the $x=0.3$ single crystal reveals a discontinuous change in the scattering at the Curie temperature for wave vectors below $\approx 0.065 \text{ \AA}^{-1}$. Strong relaxation effects are observed for this domain scattering, for the magnetic order parameter, and for the quasielastic scattering, demonstrating that the transition is not continuous in nature, in good agreement with the temperature dependence of the central component of the magnetic fluctuation spectrum, the polaron correlations, and the spin stiffness reported previously. This behavior contrasts with the continuous behavior observed for the $x=0.2$ crystal, which is well away from optimal doping. There is a large oxygen isotope effect observed for the T_C in the polycrystalline samples, and the Curie temperature is decreased by 7 K by substituting 50% ¹⁸O in the $x=0.33$ sample. For the optimally doped $x=3/8$ sample we observed $T_C(^{16}\text{O})=266.5 \text{ K}$ and $T_C(^{18}\text{O})=261.5 \text{ K}$ at 90% ¹⁸O substitution. Although T_C is decreased by 5 K for the $x=3/8$ sample the temperature dependence of the spin-wave stiffness is found to be identical for the two samples. These results indicate that T_C is not solely determined by the magnetic subsystem, but instead the ferromagnetic phase is truncated by the formation of polarons which cause an abrupt transition to the paramagnetic, insulating state. The application of uniaxial stress in the $x=0.3$ single crystal sharply enhances the polaron scattering at room temperature. Measurements of the phonon density-of-states show only modest differences above and below T_C and between the two different isotopic samples.

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I. INTRODUCTION

One of the simplifying features of conventional isotropic ferromagnets such as Fe, Co, Ni, and EuO is the isolation of the spin system from the underlying lattice. This made these types of materials ideal candidates to investigate the critical dynamics, scaling behavior, and nature of second-order phase transitions. This situation contrasts dramatically with the colossal magnetoresistive (CMR) oxides, exemplified by La_{1-x}Ca_xMnO₃, where the electronic, lattice, and magnetic degrees of freedom are intimately intertwined.¹⁻⁸ This coupling leads to a number of interesting physical phenomena, including orbital ordering, charge ordering, polarons, stripes, ferromagnetic droplet formation at small^{9,10} and large x , and the dramatic decrease in the electrical resistance in the ferromagnetic-metallic state that is the CMR effect. Investigations of these phenomena and the close relationship between CMR compounds, relaxor ferroelectrics, and high temperature superconductors have made this a field of intense study. An additional strong motivation comes from the possibility of technical applications of CMR materials in a new generation of magnetic sensors, read/write heads, and for a variety of new spin-sensitive electronics.

In the CMR materials the transport and magnetic order are connected through the ferromagnetic double exchange mechanism¹¹ while the Jahn-Teller effect locally distorts the lattice that surrounds a Mn³⁺ ion,^{1-3,8,12-15} coupling the electronic and magnetic degrees of freedom with the lattice. The

resulting ferromagnetic transition was found to be dramatically different from conventional isotropic ferromagnets, and was interpreted as a discontinuous transition between the ferromagnetic-metallic state and the paramagnetic-insulating state.¹⁶ In particular, the spin wave stiffness was found not to renormalize to zero at the Curie temperature as it should in a second-order transition. Instead, the spin wave intensities decreased rather than increasing, and the spectral weight of the magnetic fluctuations shifted to a quasielastic spin-diffusive peak in the fluctuation spectrum that rapidly developed as $T \rightarrow T_C$ and peaked near T_C .¹⁶⁻¹⁸ Concomitant with this spin-diffusive component is the development of diffuse lattice (polaron) scattering, that is correlated in a way that is consistent with the formation of zig-zag, CE-type structural stripes.¹⁹ The intensity for these polaron-polaron correlations also increases dramatically as $T \rightarrow T_C$ in a manner very similar to the spin-diffusion component, and the temperature dependence for both of these closely follows the resistivity of the metal-insulator transition.^{5,18,20-26} It has recently been observed that these polarons form a glass above the ferromagnetic-metallic to paramagnetic-insulator transition, which subsequently melts at a higher transition temperature to a fluid.^{27,28}

In the following we present a variety of neutron scattering measurements associated with the ferromagnetic transition in La_{1-x}Ca_xMnO₃. The focus of these measurements is to show that the ferromagnetic transition in the region from $x=0.3$ to $x=3/8$, the optimally doped regime, is first order.

In Sec. III we present neutron scattering results on a single crystal of $\text{La}_{0.8}\text{Ca}_{0.2}\text{MnO}_3$. This is below the optimal doping regime and the behavior is consistent with a conventional second-order ferromagnetic transition. In Sec. IV we investigate several samples in the region of optimal doping. In Sec. IV A we focus on the relaxation effects associated with a first-order ferromagnetic transition. These effects are observed in small angle neutron scattering (SANS), magnetic Bragg scattering, and the quasielastic spin diffusive peak in the magnetic fluctuation spectrum.^{16,17} In Sec. IV B we include some intriguing results concerning the effect of uniaxial stress on the ferromagnetic transition focusing on the polaron formation above T_C . These results show a further example for how strongly external perturbations affect the magnetic transition in CMR materials. In Sec. IV C we consider optimally doped samples specifically with $x=3/8$ (the highest T_C on the phase diagram) which have been treated with ^{16}O and ^{18}O isotopes. We study the temperature evolution of the spin wave spectrum as has been done previously for single crystals and powders with lower x values.¹⁶⁻¹⁸ The results demonstrate a ferromagnetic transition that is clearly first order in nature, as interpreted in the original measurements on this system. The results also make it clear that the observed T_C is determined not by magnetic interactions but instead by polaron formation; the formation of polarons truncates the ferromagnetic-metallic state. Finally, in Sec. IV D we present our measurements of the phonon density-of-states both above and below T_C for both isotopes of oxygen, and observe small but measurable differences in the lattice dynamics.

II. EXPERIMENTAL DETAILS

The single crystal samples of $\text{La}_{0.7}\text{Ca}_{0.3}\text{MnO}_3$ and $\text{La}_{0.8}\text{Ca}_{0.2}\text{MnO}_3$ were grown by the floating zone technique and have a mass of ≈ 0.7 g. The crystal mosaics were single peaked with an intrinsic width less than a $\frac{1}{4}$ degree. In this range of compositions the crystal structure is orthorhombic but the distortion is small and the crystallographic domains appear to be equally populated. Therefore, on average, the sample itself is cubic, and we use cubic notation for simplicity. Nearest-neighbor manganese atoms are then along the [100]-type directions. The crystals were mounted in the ($hk0$) plane and the cubic lattice parameter was $a=3.867$ Å at room temperature. Elastic and inelastic neutron scattering measurements were collected on the BT-2 and BT-9 thermal triple axis spectrometers at the NIST Center for Neutron Research, using pyrolytic graphite monochromators, analyzers, and filters. Small angle neutron scattering data were obtained on the NG-7 30 m SANS spectrometer.

The $x=0.3$ single crystal has a Curie temperature of 251 K, which is the highest reported T_C for any Ca doped single crystal. This is not the highest possible T_C but it has not proved possible yet to grow single crystals with a higher Ca content and therefore bulk studies at higher x are restricted to polycrystalline samples. For the isotope studies a 20 g polycrystalline sample of $\text{La}_{5/8}\text{Ca}_{3/8}\text{MnO}_3$ was specifically prepared for the measurements, using the standard solid-state reaction technique. This Ca concentration was

chosen because it is at the maximum T_C and therefore any change in the Curie temperature cannot be caused by small changes in the overall oxygen content.^{29,30} At each stage of the measurements high resolution powder diffraction measurements were carried out using the BT-1 powder diffractometer to characterize the powder sample(s), and at each stage the samples were found to be single phase. To implement the isotope substitution, the original sample was divided in half and each half was heat treated in parallel with either ^{16}O and ^{18}O , respectively, through an established procedure.¹² The fraction of isotopic substitution was determined by measuring the weight change, which established that approximately 90% substitution of ^{18}O was achieved. Bulk magnetization measurements were taken using a commercial SQUID magnetometer,³¹ while the ferromagnetic contribution to the (100) Bragg peak was measured by neutron diffraction (data presented in Sec. IV C). The nuclear scattering amplitudes for ^{16}O and ^{18}O are indistinguishable so that the total oxygen content can be determined by diffraction, but the fraction of substitution cannot be determined by this technique. After the ^{18}O inelastic and diffraction measurements were completed, the sample was converted back to ^{16}O and the properties were re-measured to establish that there was no change in the composition or changes in properties from the initial sample, to within experimental error. Finally, measurements were also carried out on the original $x=0.33$ powder sample^{16,32} for comparison purposes. This sample was re-heat treated at higher temperatures compared to the original preparation (and a small amount of Ca may have been lost in this process), and then the isotope substitution was performed. In these initial measurements only 50% substitution was achieved due to the limited amount of ^{18}O available at the time, but the shift in T_C was still larger than at optimal doping.

BT-2 and BT-9 thermal triple-axis spectrometers were employed to measure the magnetic Bragg scattering and determine the spin dynamics. For doping x between $x \approx 0.15$ and $x \approx 0.5$ $\text{La}_{1-x}\text{Ca}_x\text{MnO}_3$ compounds are ferromagnetic at low temperatures.^{29,30} At smaller x there is anisotropy in the dispersion relations and additional excitations,³³ but in the regime of interest here the system behaves to a very good approximation as an ideal isotropic ferromagnet. In polycrystalline samples this makes it possible to measure the spin dynamics around the forward scattering direction using the (000) peak reciprocal lattice point,^{16,34} while of course in single crystals the dispersion relations and linewidths can be measured around any reciprocal lattice point.³⁵ In either case the long wavelength excitations in a ferromagnet should have quadratic dispersion,

$$E_{SW} = Dq^2 + \Delta, \quad (1)$$

with a temperature-dependent spin wave stiffness D and a possible spin gap Δ that would arise from anisotropy. The spin wave spectrum has been found to be gapless within the precision of neutron inelastic scattering measurements,⁵ which indicates that these materials are to a very good approximation rotationally isotropic.

The phonon density-of-states was measured on the Filter Analyzer Neutron Spectrometer (FANS) and the time-of-

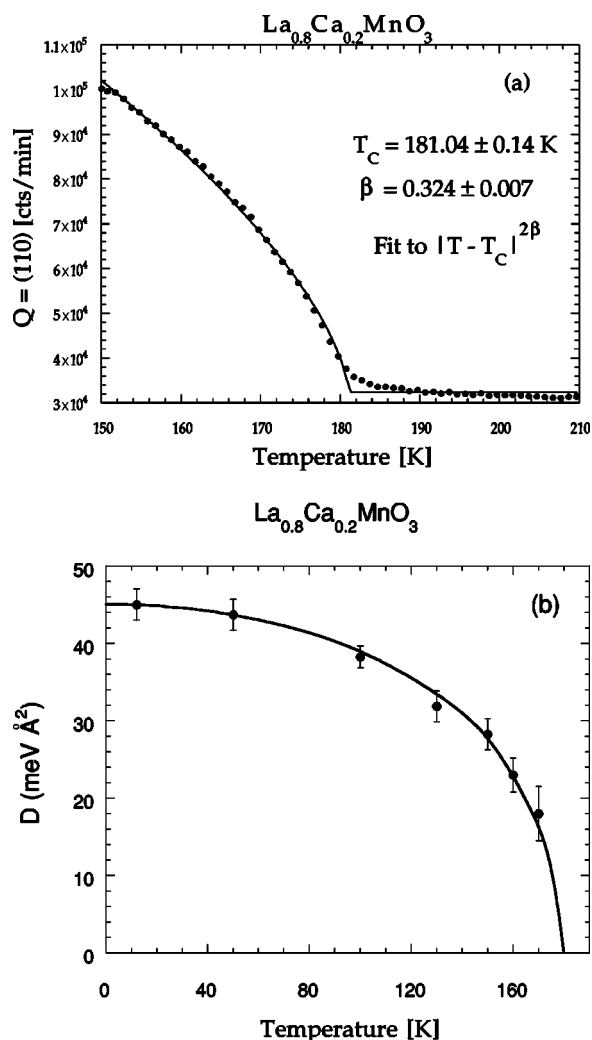


FIG. 1. (a) Measurement of the ferromagnetic order parameter at the (110) peak for a single crystal of $\text{La}_{0.8}\text{Ca}_{0.2}\text{MnO}_3$ near T_C . Data are well described by a power law with an exponent that agrees with a three-dimensional Heisenberg model. (b) Spin wave stiffness for the same sample. D renormalizes to zero as $T \rightarrow T_C$, as expected for a conventional ferromagnetic transition that is continuous in nature.

flight Fermi Chopper Spectrometer (FCS). Uncertainties quoted for all the neutron measurements are statistical and represent one standard deviation.

III. MAGNETIC PROPERTIES FOR $x=0.2$

In Fig. 1(a) we show a measurement of the ferromagnetic (110) Bragg peak in the vicinity of the Curie temperature for the $x=0.2$ single crystal. Above the Curie temperature only the Bragg peak due to nuclear scattering is observed, while below T_C the magnetic contribution to the peak intensity is proportional to square of the order parameter (magnetization). This magnetic intensity is well described by a power law, except in the immediate vicinity of the Curie point where critical scattering is significant. The data above 180 K were therefore excluded from the fit. The fit gives $T_C = 181.04 \pm 0.14 \text{ K}$, and an (effective) exponent β

$= 0.324 \pm 0.007$ which is close to the prediction of the three-dimensional Heisenberg model ($\beta=0.365$).³⁶ The errors represent statistical uncertainties only, and no account was taken of a possible spread of transition temperatures since the critical properties are not central to the present investigation. This Curie temperature is in good agreement with the values of 175 K,³⁷ 178 K,³⁸ and 185 K³³ reported previously for single crystals of the same nominal composition.

The spin wave spectrum has been measured at small wave vectors, and the spin wave stiffness coefficient D is shown in Fig. 1(b) as a function of temperature. Each constant- q scan was least-squares fit to a quadratic dispersion law convoluted with the instrumental resolution. Correcting for the spectrometer resolution has the effect of decreasing the observed value of D by $\sim 5-10\%$, and the low temperature value was determined to be $46 \pm 2 \text{ meV \AA}^2$ for this single crystal. This is in good agreement with the values previously reported for this composition.^{33,37,38} As the Curie temperature is approached the spin wave intensities were found to increase due to the increase in the thermal population of spin waves, caused by the combined higher temperatures and lower $D(T)$ values. In Fig. 1(b) we see that D appears to renormalize to zero at T_C , as expected for a conventional second-order ferromagnetic transition. We did not find the development of a central quasielastic component to the magnetic fluctuation spectrum, in contrast to the behavior found in the metallic regime.^{16-18,38} For the $x=0.20$ composition, Dai *et al.*³⁸ also did not find a central component, but they did suggest that $D(T_C)$ was finite, which we do not see in the present measurements. However, we did not measure the spin wave stiffness quite as close to T_C as Dai *et al.*³⁸ (our final value is taken 10 K below T_C and theirs is 5 K below T_C) so our measurements cannot refute their result (which also states that D fails to renormalize at $x=0.25$). It is reported by Biotteau *et al.*³³ that D renormalizes to zero for the $x=0.125$ composition so there is some definite change in the renormalization behavior of D between $x=0.125$ and $x=0.25$.

IV. MAGNETIC PROPERTIES NEAR OPTIMAL DOPING

The overall behavior changes as x increases to the optimal doping range for ferromagnetism. The excitations at long wavelengths are still isotropic spin waves with quadratic dispersion, but as mentioned in the Introduction D fails to renormalize to zero at T_C . Instead a quasielastic peak appears in the fluctuation spectrum, indicative of the type of spin diffusion that occurs above T_C . This was interpreted as a coexistence between the low- T ferromagnetic-metallic phase, and the high- T paramagnetic-insulator phase, with the phase fractions changing with temperature.^{16,39} The spin diffusion component then grows at the expense of the spin wave intensities. Measurements of the magnetic correlation length as one approaches T_C from above show that this length does not diverge at the transition as would be expected for a continuous transition.^{16,18,40} This agrees with measurements of the paramagnetic scattering, and the correlation length is only weakly temperature dependent at $\sim 15 \text{ \AA}$ through the transition. These features cannot be explained in the context of a

conventional second-order ferromagnetic phase transition, but can be understood if one assumes that long range ferromagnetic order is not the sole order parameter of the transition and that the coupling to the other order parameter forces the transition to be first order.⁴¹

A. Relaxation effects near T_C

For a conventional ferromagnet above T_C the magnetic scattering is concentrated around the Bragg peaks with q -dependent intensity that follows the Ornstein-Zernike form

$$I \propto \frac{1}{q^2 + \kappa^2}. \quad (2)$$

The correlation length $\xi (=1/\kappa)$ in real space is the parameter that diverges at T_C for a second-order transition. Such a divergence was not observed in previous polycrystalline measurements,^{16,40} or in more recent single crystal measurements made around Bragg peaks.^{17,18} However, SANS measurements on polycrystalline samples are restricted due to the large amount of small angle metallurgical scattering from the crystallites, and therefore we have pursued SANS measurements on the $x=0.3$ single crystal where high quality data could be obtained over a wider range of wave vectors and temperature, and where we have also been able to measure the time-dependence of the scattering.

The results are quite different from conventional second-order behavior, as (now) expected. The Ornstein-Zernike behavior in Eq. (2) is maintained for $q \gtrsim 0.065 \text{ \AA}^{-1}$ with a short and weakly temperature-dependent correlation length as found previously.^{16,18} In Fig. 2(a) we show data at smaller q above and below T_C . There is a discontinuous increase in the scattering at T_C which cannot be described in terms of the Ornstein-Zernike form or the common Porod form ($\log I \propto \log q$ with a slope of -4).⁴²⁻⁴⁴ This scattering is strongly peaked in the forward direction with a slope of -5 on a Porod plot,⁴³ and dominates the intensity on the (two-dimensional position-sensitive) detector. Figure 2(b) shows the temperature variation of this scattering, recorded as the total counts on the detector. This corresponds to an integration of the scattered intensity from roughly $q = 0.001 \text{ \AA}^{-1}$ to $q = 0.015 \text{ \AA}^{-1}$. The wavelength is equal to 10 \AA and the detector was 15 m from the sample for this measurement. Without the metallurgical details of the crystal (crystallographic type and density of twin domains and defects, for example) a more quantitative analysis is unwarranted. The abrupt nature of the onset of the scattering shown in Fig. 2(b) is consistent with the transition being discontinuous.

One hallmark of first-order transitions is that, even if thermal hysteresis is too small to be observed, they are accompanied by relaxation effects. Figure 3 shows SANS measurements at three different wave vectors as a function of time, after warming to 250 K . The time for the sample to equilibrate near T_C is clearly quite long. The data can be described by an exponential decay with relaxation time τ , and the dashed curve is a fit of this form to the data. Different values of q gave similar values of τ within statistical uncertainties, showing that all length scales measured in these SANS ex-

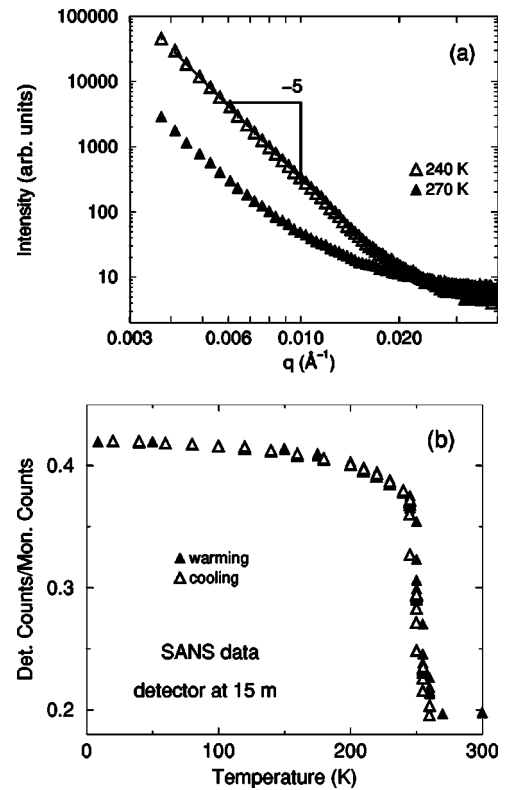


FIG. 2. (a) The q dependence of the small angle neutron scattering intensity above and below T_C for a single crystal of $\text{La}_{0.7}\text{Ca}_{0.3}\text{MnO}_3$. Below T_C there is a significant increase in scattering concentrated in the forward direction with a slope of -5 on a log-log scale. (b) Integrated intensity as a function of temperature showing the abrupt change at $T_C = 251 \text{ K}$.

periments exhibit similar relaxation times. Warming was performed in steps of 5 K every 40 minutes for these measurements, and the relaxation time was found to be maximized at T_C .

One of the difficulties in unambiguously identifying this as a first-order transition has been that, although D does not renormalize to zero, the magnetic contribution to the intensity of the Bragg peaks appears to be continuous. In particu-

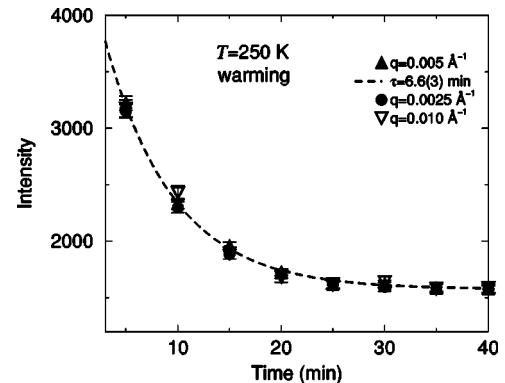


FIG. 3. Time dependence of the observed SANS scattering for three different wave vectors at 250 K . Data are scaled at $t=0$ and shifted vertically to the same asymptotic value, and are fit to a decaying exponential with a time constant $\tau = 6.6 \pm 0.3$ minutes.

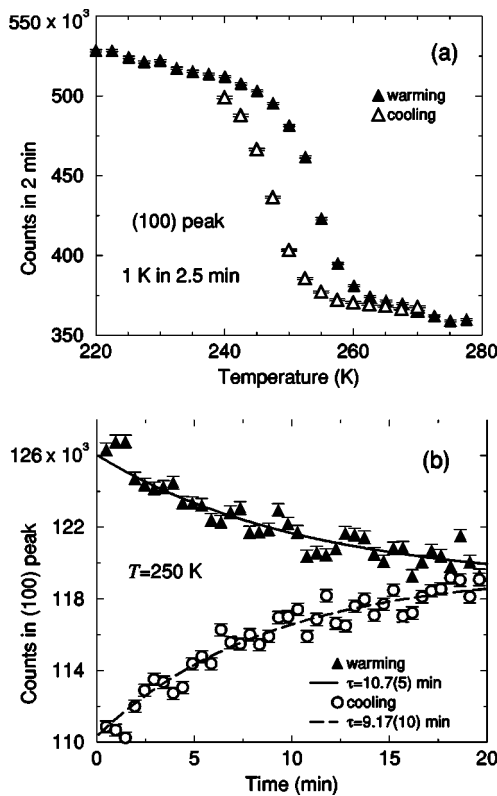


FIG. 4. (a) Thermal irreversibility of the (100) magnetic Bragg peak on warming and cooling. The data were obtained by changing the temperature in 1 K steps every 2.5 min. (b) Time dependence of the intensity of the (100) Bragg peak, both warming to the set point of 250 K, and cooling to the same set point. Data are fit to an exponential relaxation. Note that the time constants are roughly identical and the asymptotic values are identical within experimental uncertainties, so there is no genuine hysteresis observable.

lar, although there are some “irreversibilities” observed on warming and cooling,¹⁶ no clear hysteresis has been identified. The availability of the $x=0.30$ single crystal has allowed us to investigate the behavior of the magnetic Bragg peak in detail. In Fig. 4(a) we plot the intensity of the (100) Bragg peak, and we see a clear difference on warming and cooling. This apparent hysteresis is due to relaxation effects, and in Fig. 4(b) we show the time dependence of the scattering, after the sample temperature arrives at the set point. The curves are fits to exponential relaxation functions, which describe the data quite well. The fitted τ versus temperature is shown in Fig. 5, on warming and cooling. Although the statistical uncertainties are large, particularly when the total change in scattering between two temperatures is small, we can clearly see that τ is maximized near T_C .

A further result of the time-dependent measurements is that the asymptotic value of the magnetic contribution to the Bragg peak at any particular temperature is identical on warming and cooling, as shown in Fig. 6. Thus there is no true thermal hysteresis in the asymptotic time region. We note that the value for T_C determined in these measurements, 251 ± 3 K, is slightly lower than the preliminary value we reported previously 257 ± 1 K.¹⁸ The (correct) value of 251 K is determined here by locating the inflection point,

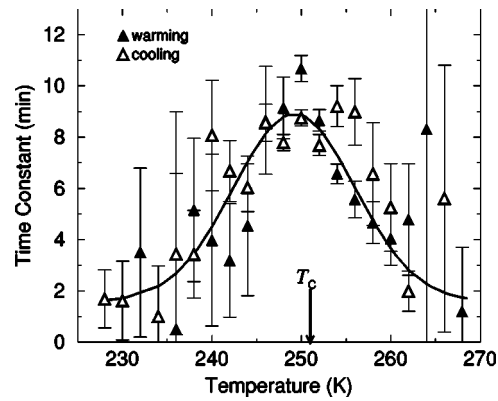


FIG. 5. Relaxation time τ as a function of temperature for the (100) Bragg peak. The short time constant found away from the transition is typical of the time it takes for a sample to equilibrate, but in the vicinity of T_C the time constant increases dramatically, indicating that the transition is first order in nature.

after taking the data on warming and cooling and properly taking into account the long time-scale behavior of the system. This value is also in good agreement with the abrupt onset of SANS shown in Fig. 2(b). The SANS data also lack true thermal hysteresis.

We also measured the relaxation of the central quasielastic component of the magnetic fluctuation spectrum just below T_C for the $x=3/8$ ^{16}O sample, and the data are shown in Fig. 7. (This sample will be described in more detail in Sec. IV C.) The sample temperature was changed to 266 K from well below or well above this temperature, and the data collection at the elastic position was started when the sample block reached the set point. The measured relaxation times are somewhat different for the two cases, but it is clear that the asymptotic values are identical to within statistical uncertainties. Hence, the activation barrier and latent heat for this transition are rather small, which has hindered the identification of this as a first-order transition in the specific heat until recently.⁴⁵

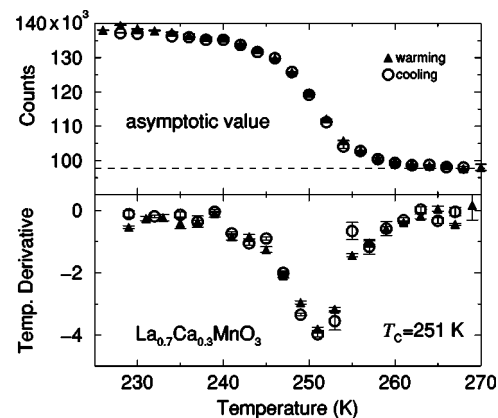


FIG. 6. Intensity of the (100) magnetic Bragg peak, in the long-time limit. The data reveal that there is no genuine hysteresis observable. The bottom curve shows the derivative of the data, where the minimum identifies a Curie temperature of 251 K.

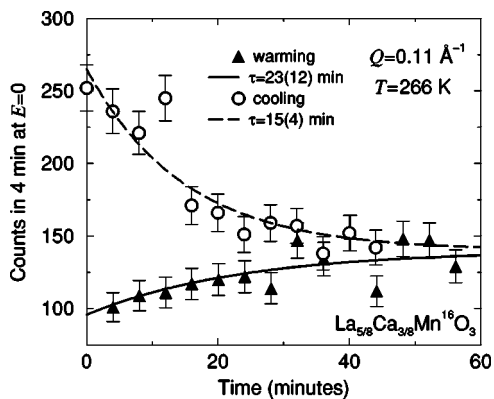


FIG. 7. Intensity of the quasielastic magnetic scattering for a temperature of 266 K at a wave vector of 0.11 \AA^{-1} . Note that it takes about an hour for the intensities to become equal on approaching from low versus high temperatures. In the long-time limit there is no genuine hysteresis observable.

B. Effect of uniaxial stress

One of the characteristic features of CMR materials is their extraordinary sensitivity to external perturbations. The application of a magnetic field, for example, represents a very modest energy in a system that orders at 250 K, but a magnetic field can drive the metal-insulator transition and the associated CMR effect. Stress, which is always present in thin films, also has a dramatic effect on the electronic properties, and is one of the primary reasons for the rejuvenation of interest in this class of materials.⁴⁶ We have already established that the formation of polarons in the paramagnetic state is directly related to the resistivity,¹⁸ and we may then anticipate that the polarons themselves will be directly affected by stress.

To date there have been only a few studies of CMR as a function of applied force, using two-dimensional strain in thin films⁴⁷ and hydrostatic pressure.^{37,48} In addition, the effects of internal pressure that result from the substitution of smaller Pr ions for the La A-site cation, for example, have been performed. Other options for cation doping have shown that the behavior of the entire x - T La-Ca phase diagram can be achieved via changing the average radius at the A-site.^{41,49,50} Results can often be understood in terms of the double exchange model when the distortion of the ideal perovskite lattice is considered. Changing the average radius of the A-site cation changes the tilt of the oxygen octahedra that surround the B-site cation, which affects the electronic bandwidth and the double exchange. An increase in distortion away from the ideal perovskite structure tends to favor charge-ordered insulating behavior, where the antiferromagnetic interaction dominates over metallic ferromagnetism.

Our single crystal of $\text{La}_{0.7}\text{Ca}_{0.3}\text{MnO}_3$ was mounted in a stress rig on the BT-9 triple axis spectrometer in the $(hk0)$ zone. Stress was applied along the axis of the crystal, roughly 30 degrees from (001) in the direction of (220). All measurements were made at room temperature where the charge peak intensity has dropped to roughly 30% of its maximum value near T_C . Given the geometric limitations imposed by the stress rig we were essentially limited to mea-

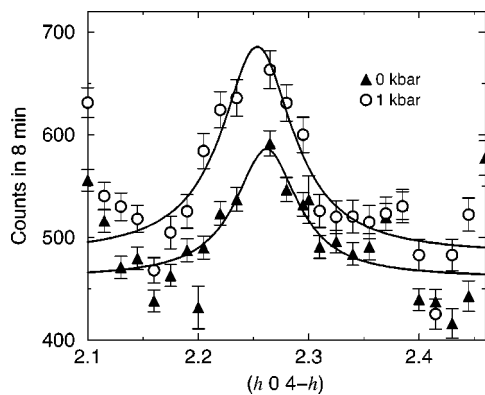


FIG. 8. Change in the intensity of the charge-order peak when uniaxial stress is applied. The intensity increases by more than a factor of two in going from 0 kbar to 1 kbar of uniaxial stress. The intensity of the single-polaron (diffuse) scattering also increases with applied stress.

surements near the (220) peak using an incident energy of 52 meV. The charge order peaks are broad (corresponding to short range correlations) and occur at wave vectors $(\pm\frac{1}{4}, \pm\frac{1}{4}, 0)$. Near (220) these occur at (1.75 2.25 0) and (2.25 1.75 0); the structure factor is negligible for the “longitudinal” peaks at (1.75 1.75 0) and (2.25 2.25 0). In Fig. 8 we show a transverse scan of the charge order peak at (2.25 1.75 0) at 0 kbar and 1 kbar. The intensity of the peak has increased by nearly a factor of two, while the intensity of the fundamental Bragg peaks are unaffected at this level of stress. It is evident that the diffuse scattering, which is related to the “single-polaron” diffuse scattering, has also increased. Uniaxial stress therefore has a rather dramatic effect on the polaron scattering.

Because of the rather large change in intensity that was observed when 1 kbar of stress was applied, we reduced the force to zero in order to determine the behavior at lower values of the stress. The polaron intensity was reduced, but not to the initial value; the observed scattering was about halfway between the initial data and the 1 kbar data. Subsequent measurements up to 0.5 kbar showed no change, and with a further increase the scattering began to increase up to the 1 kbar data. Hence we believe that the scattering would increase approximately linearly with stress from the virgin state. Our current measurements show that cycling up to 1 kbar of stress has some residual effects, but without further measurements we cannot say if this affects T_C or can be seen by a full measurement of the residual stress. Unfortunately, shortly after returning to 1 kbar of stress, the crystal shattered, ending the data collection.

C. Oxygen isotope effects

One very strong indication for some type of coupled transition in $\text{La}_{1-x}\text{Ca}_x\text{MnO}_3$ is the oxygen isotope effect.^{12,51-57} since specific isotopic masses of the components do not usually affect the magnetic interactions. However, in $\text{La}_{0.8}\text{Ca}_{0.2}\text{MnO}_3$ substitution of ^{18}O for ^{16}O leads to a 20 K decrease in T_C , from 180 K to 160 K.¹² An isotope effect in antiferromagnetic $\text{La}_{0.5}\text{Ca}_{0.5}\text{MnO}_3$ is also seen, but in the op-

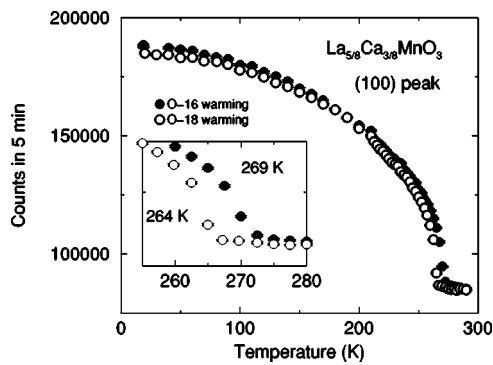


FIG. 9. Measurement of the intensity of the magnetic scattering on warming for polycrystalline samples of $\text{La}_{5/8}\text{Ca}_{3/8}\text{MnO}_3$ that have been treated with ^{16}O and ^{18}O . The Curie temperature is reduced by 5 K for the ^{18}O sample.

posite direction and affecting the Néel temperature.⁵¹ To investigate the effect of oxygen isotope substitution on the magnetic behavior we chose to work specifically at $x=3/8$, where T_C is maximized and the CMR effect is large.³⁰ This choice of x minimizes any possible variation in or distribution of T_C with the overall oxygen content. Indeed, for the measurements in Ref. 12 the observed oxygen isotope effect has been partially attributed to different oxygen stoichiometry.⁵³ For the present samples the high resolution powder diffraction measurements could be used to establish that the oxygen contents for both samples were identical to within experimental uncertainties, and thus there can be no significant variation in T_C due to the total oxygen content. The lattice parameters, average atomic positions and occupancies were identical as well. Density measurements indicate that roughly 90% of the ^{16}O was replaced by ^{18}O in the substituted sample.

Figure 9 shows the intensity of the (100) powder peak as a function of temperature for the two samples, to determine the effect of oxygen isotope substitution on T_C . The data were obtained as a function of warming at identical heating rates. The Curie temperature of the ^{18}O sample is 5 K lower than the ^{16}O sample, as can be seen in the inset. The same 5 K difference was seen on cooling, along with a 5 K irreversibility in both the curves relative to the warming runs. Given the first-order nature of the transition the irreversibility is expected and is consistent with previous measurements. Taking the average of these warming and cooling runs gives $T_C(^{18}\text{O})=261.5\pm 0.5$ K and $T_C(^{16}\text{O})=266.5\pm 0.5$ K. As we pointed out earlier for the single crystal measurements, these T_C 's are determined by taking the inflection points of the intensity of the magnetic scattering curves. The value of T_C for the ^{16}O sample is indeed the maximum value seen on the Ca-concentration phase diagram and is a further indication of the high quality of these samples.

This 5 K shift is much smaller than reported in the $x=0.2$ sample, so for comparison we have also made measurements on the original $x=0.33$ sample.¹⁶ With only 50% ^{18}O substitution we obtained $\Delta T_C=-7$ K. It is noteworthy that the transition for the partially substituted sample was just as sharp as for the fully substituted sample, indicating that there is no significant effect on the properties caused by the oxy-

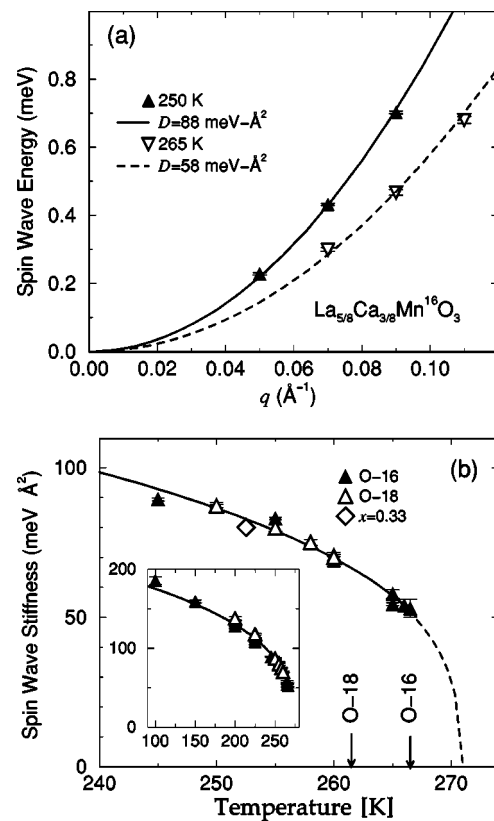


FIG. 10. (a) An example of the spin wave dispersion relation for two typical temperatures. Data are well described by gapless quadratic dispersion with the indicated values of D . (b) D versus temperature for the ^{16}O and ^{18}O samples. The observed D values are identical at each temperature independent of the oxygen isotope, despite the difference in T_C (as indicated by the arrows). The extrapolated temperature where $D(T)$ would renormalize to zero is 271 K, but the ferromagnetic-metallic state abruptly disappears at the temperature where the polarons form, and this truncation occurs at a lower temperature for the ^{18}O sample. The $D(T_C)$ value for the original data in Ref. 16 with $x=0.33$ is also indicated by the open diamond symbol. The inset shows the stiffness over a wider temperature range.

gen isotopic randomness. The larger isotopic shift with decreasing x that we find is in good agreement with the trends already established for this system.

Having confirmed that oxygen isotope substitution changes T_C , the next question is how the magnetic interaction strength changes, which we characterize by measuring the spin dynamics. In Fig. 10(a) we show the low- q spin wave dispersion in the ^{16}O sample at 250 K and 265 K. The dispersion is well described by Eq. (1) with gapless excitations, typical of these isotropic ferromagnets.^{16-18,24,58} The spin wave stiffness has decreased considerably over this temperature range upon approaching T_C . Similar measurements were made to determine D in the ^{18}O sample. Figure 10(b) shows a comparison of D versus T for the two samples. The two samples show *identical* results for D over the entire temperature range where spin waves are observed, i.e., below the respective Curie temperatures. Note also that there are in fact two complete sets of data for the ^{16}O : one taken with the two samples heat treated in parallel, and one set taken after

exchanging the ^{18}O for ^{16}O . Both sets of data are identical. The only difference found between the samples with the different oxygen isotopes is that for the lower- T_C ^{18}O sample the measurements cannot be continued to as high a temperature as with the ^{16}O sample. This result obviously cannot be understood in the conventional context of ferromagnetism since T_C has changed but the ferromagnetic interaction as measured by D has not. Hence it is clear that T_C is not determined by magnetic interactions alone.

The $D(T)$ data can be well described by a power law in this temperature range, with an extrapolated value of $T_C = 271$ K which is identical for both isotopes. The exponent used in this fit is $\beta - \nu = 0.4$, which is fairly close to the prediction of the Heisenberg model $\beta - \nu = 0.324$.³⁶ The particular value of the exponent is not the central issue here; the point is that the exchange interaction is not affected by the isotope substitution, and the extrapolated T_C is identical for both samples. For both samples the actual magnetic transition occurs before D has renormalized to zero, but note that the renormalization has progressed much farther than is the case for samples with lower x , lower T_C samples. Indeed, Fernandez-Baca *et al.*¹⁷ have established that D changes little over a wide range of tolerance factor, while large changes in T_C occur. From this point of view the (O isotope-independent) magnetic interactions determine this extrapolated T_C , but polaron formation truncates the transition, thereby determining the observed T_C . This truncation occurs in a first-order manner and leaves D at a finite value, and is the determining factor for T_C over a range of x . Other measurements of $D(T)$ for smaller $x \sim 0.3$ compounds are very similar to the $x = 3/8$ samples, just with a different truncation value at the different T_C 's. The downturn in T_C with decreasing x in the phase diagram is then caused by polaron formation rather than a weakening of the double exchange mechanism. This behavior continues until the system becomes insulating, and $D(T=0$ K) suddenly drops. For our insulating $x = 0.2$ sample $D(T=0$ K) is 3 times smaller than in the optimally doped sample, in good agreement with the results of Refs. 38 and 59.

This rapid truncation of ferromagnetism is clearly revealed in Fig. 11, which shows the magnetic fluctuation spectrum as a function of temperature at a wave vector $q = 0.11 \text{ \AA}^{-1}$ for the two samples. They show a temperature evolution similar to what has been observed for previous measurements of samples in the optimally doped regime as described at the beginning of Sec. IV, the main difference being that the change in the spin wave spectrum with temperature is very abrupt. The elastic incoherent nuclear scattering and a flat temperature-independent background have been removed from the spectra in Fig. 11, and the curves are fits to the spin wave dispersion [Eq. (1)] plus a quasielastic (spin diffusion) spectral weight function convoluted with the instrumental resolution. Spectra at a series of q 's were obtained at each temperature, and the data collection time between runs at different temperatures is sufficiently long that there are no irreversibility effects due to warming and cooling. For the ^{16}O sample [Fig. 11(a)] we see a three component spectrum at 265 K, just below T_C . The spin waves are observed in energy gain ($E < 0$) and energy loss ($E > 0$) at

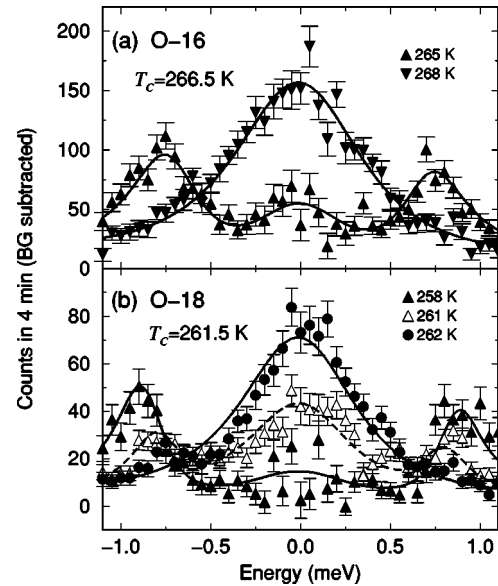


FIG. 11. Constant- q scans at a wave vector of 0.11 \AA^{-1} showing the nature of the magnetic fluctuation spectrum in the vicinity of T_C for the ^{16}O and ^{18}O samples. (a) For the ^{16}O sample there is a three peak structure at 265 K, while just above T_C at 268 K there is a single quasielastic peak. (b) For the ^{18}O sample at 258 K most of the weight of the magnetic scattering is in the spin wave excitations. At 261 K the weight in the spin wave scattering has decreased, and the spectral weight has shifted to the central component. At 262 K the scattering is completely diffusive in nature. There is no evidence of the usual collapse of the spin wave energies found in a second-order transition. Note that the count rates for the two samples are different even though the samples were approximately the same size, and this difference is due to the data being collected under different experimental conditions; ^{16}O data were taken on BT-2 with a collimation of $20' - 10' - 10' - 20'$, while the ^{18}O data were collected on BT-9 with a collimation of $10' - 13' - 10' - 16'$.

~ 0.8 meV. At 268 K, just above T_C for this sample, the spin waves have disappeared and the scattering is purely quasielastic (peaked at $E=0$). For the ^{18}O sample [Fig. 11(b)] we include one scan a bit farther below its respective T_C where we observe well defined spin waves at 258 K along with a small quasielastic component fitted to the data. At 261 K the central component has increased to form a clear peak, while the spin waves have lost over 50% of their intensity and have moved to lower energies, consistent with all previous observations of the lower x compounds which show growth of the central peak at the expense of the spin waves. But further increasing temperature by only 1 K wipes out the spin waves entirely, leaving only a central peak that corresponds to ferromagnetic correlations in the paramagnetic phase. This very rapid temperature evolution is evidence for the truncation of the magnetically ordered phase in a first-order manner. The nature of the magnetic fluctuation spectrum is the same for both samples, except that the transition for the ^{16}O sample is extended higher in temperature by 5 K. This increased T_C means that the spin wave stiffness at their respective Curie temperature $D(T_C)$ is about 30% lower than for the ^{18}O sample compared to the ^{16}O sample. Note in particular that the data for the ^{18}O sample is pure spin diffusion for 262 K,

while for the ^{16}O sample we have propagating excitations at 265 K. We remark that we have taken data for temperatures between the ones shown in Fig. 11 to examine the possibility that there is a sudden decrease in the spin wave energies that might be attributable to a sudden decrease in the effective exchange energy, but all the data indicate that we have two-phase coexistence,¹⁶ a ferromagnetic-metallic phase and a paramagnetic-polaron glass insulator. For lower x samples the truncation occurs at lower temperature, with a corresponding higher value of $D(T_C)$ as indicated in Fig. 10(b). Thus the transition at lower x is more strongly first order. We have also found that the central peak usually was observable over a wider range of temperature at smaller x , 15 K or more, compared to the ~ 5 K range observed here. This may also be related to how close the system is to the extrapolated second-order phase transition temperature T_C .^{45,58}

D. Phonon density-of-states

The phonon density-of states was measured for the $x = 3/8$ powder samples discussed in Sec. IV C using time-of-flight and filter analyzer spectrometer techniques. The time-of-flight Fermi chopper spectrometer uses cold monochromatic neutrons incident on the sample, and then detects the inelastically scattered neutrons in a bank of detectors. Near room temperature one can reliably measure energy transfers up to ~ 30 meV. The density-of-states is obtained by removing the thermal population factor, weighting the scattering by q^2 (assumes phonon scattering), and then averaging over a range of q . Note that for low energies and modest values of q a considerable part of the spectrum is magnetic in origin for these samples. We have made no attempt to separate out the magnetic contribution from the lattice (nuclear) contribution.

The filter analyzer spectrometer (FANS) uses cooled polycrystalline beryllium and polycrystalline graphite as an analyzer, so that only neutrons with energies below 1.7 meV reach the detectors. The detectors span a large solid angle (1.3 steradians), and the incident energy is scanned to determine the inelastic scattering. In contrast to the Fermi chopper spectrometer this instrument is designed to operate in an energy loss mode and no correction needs to be made for q^2 since a $1/q^2$ term enters into the resolution volume. A run with an empty sample can was used to remove background. We used a pyrolytic graphite monochromator over the range of 15 meV to 40 meV and a Cu(220) monochromator for energies between 38 meV and 120 meV. Data sets for each sample were merged together after background subtraction. A similar merging was performed to include the low-energy time-of-flight data. After merging, the data were corrected for multiphonon effects, which tended to be significant only above 75 meV. Additional experimental details are given in Ref. 60.

Results for the phonon density-of-states taken above and below T_C appear in Fig. 12. Identifiable peaks appear in the data at 17, 25, 42, 51, and 63 meV. There are no phonons above 85 meV after multiphonon corrections. These observations are in agreement with recent single-crystal measurements of the phonon spectra of $\text{La}_{0.7}\text{Ca}_{0.3}\text{MnO}_3$, which show

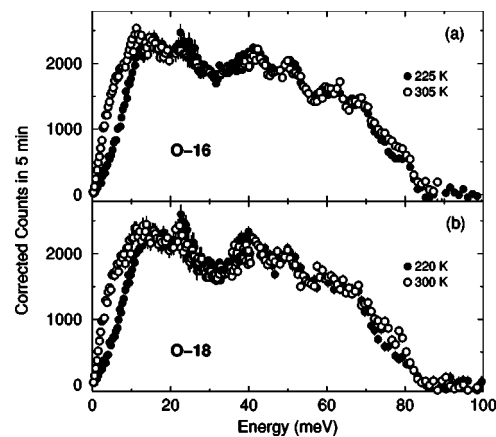


FIG. 12. Phonon density-of-states for (a) ^{16}O and (b) ^{18}O samples taken above and below T_C . Units are those appropriate for the FANS spectrometer with a Cu(220) monochromator.

no phonons above 72 meV and substantial flat regions in the dispersion branches at 23 meV and 50 meV and an optical band near 40 meV.⁶¹ Corresponding peaks in Raman scattering for an $x=0.33$ sample are found at 17, 25, and 51 meV.⁶² Other Raman peaks are also observed at 8, 10, 30, and 54 meV with a broad maximum at 85 meV. Single crystal phonon dispersion measurements on the related CMR compound $\text{La}_{0.7}\text{Sr}_{0.3}\text{MnO}_3$ show phonon branches up to 72 meV with no gaps between branches when averaging over the entire zone.⁶³ Substantial magnon-phonon interactions have also been reported,⁶⁴ but it is likely that these effects can only be observed in single crystals. For the present density-of-states data at low energies there are substantial differences above and below T_C , but this originates from the change in the magnetic scattering from spin diffusion above T_C to propagating spin waves below. At higher energies the biggest difference is in the ^{16}O data between 75 and 80 meV, where the signal at 305 K is slightly higher than at 225 K. A similar effect can be seen in the ^{18}O data except that the scatter in the data is larger. In other perovskites this energy region corresponds to the oxygen breathing modes but there was no specific peak seen in our data that could be identified with a mode of vibration in this energy region. There may be a slight decrease in intensity of the peak near 51 meV in the ^{16}O data. Both these trends are borne out by measurements at 10 K in the 40 to 100 meV range. The published single crystal results⁶¹ for the $x=0.3$ sample are far more conclusive and show significant damping effects on the longitudinal Jahn-Teller (100) breathing mode which has an energy of 72 meV at the zone center and 51 meV at the zone boundary. Near the zone boundary the intensity of the phonons strongly decreases with temperature until the mode disappears above T_C . This is consistent with a decrease in the 51 meV peak seen in our data. Similar changes are also seen in the Raman scattering.⁶² Overall, however, the changes in the phonon density-of-states are modest.

A comparison between the different samples above T_C also reveals that they are nearly identical, as shown in Fig. 13 where we compare the phonon density-of-states between samples rather than between temperatures. The different temperatures were chosen to reflect the different T_C values. The

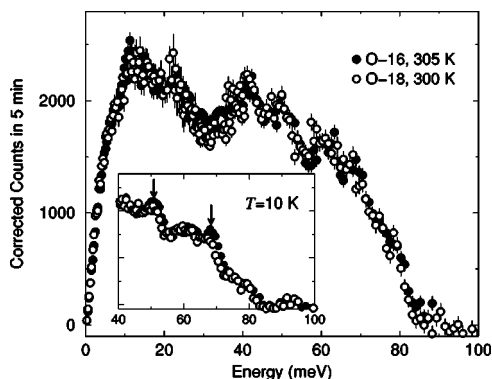


FIG. 13. A comparison of the phonon density-of-states for ^{16}O and ^{18}O samples above T_C . There are some slight differences which are more easily seen at low temperature, shown in the inset.

region that looks most affected by oxygen isotope substitution at these temperatures is between 20 and 40 meV where the scattering intensity near 30 meV is lower for the ^{18}O sample than the ^{16}O sample. Raman measurements by Franck *et al.*⁵³ near 0 K showed that the oxygen modes at 28 meV and 53 meV are shifted downward by 1 meV and 2.6 meV for the ^{18}O sample. At low temperatures (inset) there does seem to be a difference in the spectra at 68 meV where the ^{16}O peak may be slightly higher and more clearly defined. Overall, however, only small changes are observed in the density-of-states, either with a change in isotope, or with a change in temperature above versus below T_C , which is in agreement with recent tunneling data.⁶⁵ Given the role of polarons and oxygen isotope substitution in the transition this observation might be somewhat surprising. Further single crystal studies and more complete models of the transition might resolve this issue.

V. DISCUSSION

The increase in the ferromagnetic transition temperature when ^{18}O is replaced by ^{16}O , and the concomitant extension of the spin wave regime without any significant change in the overall magnetic energetics, demonstrates in an unambiguous way that the ferromagnetic state is prematurely truncated, and that the transition from the ferromagnetic-metallic state to the paramagnetic-insulator state is first order in nature. For the optimally doped $x=3/8$ composition the extrapolated Curie point, where the transition would be second order, is just above the observed ferromagnetic transition, indicating that the transition is weakly first order at this composition. It is this proximity to the second-order fixed point that enables the large ($\sim 30\%$) decrease in the value of the

spin stiffness at the respective Curie temperatures when ^{18}O is replaced by ^{16}O , making it unambiguous that the ferromagnetic state is truncated. At somewhat higher Ca concentrations the transition in fact becomes second order.⁴⁵

The origin for the first-order truncation is the sudden formation of polarons in the system, which trap the carriers and render the system insulating. It is the polaron formation temperature that determines the observed T_C in the first-order regime, not the strength of the magnetic interactions. Replacing ^{16}O by ^{18}O lowers the energy of the Jahn-Teller phonon, making it easier to form polarons and consequently lowers the transition temperature. The application of stress also lowers the barrier for polaron formation. Lowering the Ca concentration below the $x=3/8$ composition decreases the observed Curie temperature, while the exchange energy is (initially) unaffected.¹⁷ This downturn in T_C (observed) in the phase diagram then originates from lowering the energy for polaron formation, as it becomes easier for the system to form correlated Jahn-Teller distortions; recall that at $x=0$ pure LaMnO_3 is strongly Jahn-Teller distorted, forming a structure with long range orbital (and antiferromagnetic) order. The temperature range where the polaron glass is observed^{27,28} would then be expected to become larger with decreasing x . This trend continues until the system becomes insulating around $x \approx 0.2$, where the double-exchange interaction is lost and the magnetic energetics drops precipitously.^{38,59}

Ultimately it is the structure that determines the energetics for polaron formation, which can be cast in the form of the tolerance factor t .^{41,49} The Ca system has a relatively small t and forms polarons relatively easily, and the polaron formation occurs in the ferromagnetic state and causes a first-order transition. Larger ions such as Sr or Ba have a substantially larger t , and do not appear to form a significant number of polarons below T_C . The ferromagnetic transition then becomes a conventional second-order transition, that is uncoupled from the resistivity transition.^{66,67} It will be interesting to determine if there are dynamic polarons in the Sr and Ba substituted systems.

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